

On a two species chemotaxis system

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Contents

- 1-. Introduction
- 2.- One species parabolic-elliptic chemotaxis systems (known results)
- 3-. Two species chemotaxis systems, **with M. Winkler[†]**
- 4-. Two species chemotaxis system- A case of competitive exclusion,
with C. Stinner[‡] and M. Winkler[†]
- 5-. Two species chemotaxis system with nonlocal terms, **with M. Negreanu^{††}**
- 6.- The Parabolic-ODE system **with M. Negreanu^{††}**

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1 Introduction

Chemotaxis is the ability of microorganisms to respond to chemical signals by moving along the gradient of the chemical substance, either toward the higher concentration (**positive taxis**) or away from it (**negative taxis**).

Dictyostelium discoideum (Bacteria)



One species Chemotaxis systems

- Keller and Segel [1970], [1971]
- Jäger and Luckhaus [1992], Biler [1995a], [1995b], [1997]
- Herrero-Velázquez [1996], [1997], Velázquez [2002] [2004]
- T. Nagai, [1995], T. Nagai, T. Senba and K. Yoshida [1997]
- A. Friedman, A. Stevens, B. Hu, M. Mimura, Mizoguchi, Ph. Souplet Naito, Suzuki, Blanchet, Dolbeault, Perthame, Fasano Carrillo, Masmoundi, Horstmann, M. Peletier, T. Tsujikawa, K. Osaki, A. Yagi, Painter, T. Cieslak, C. Morales-Rodrigo, C. Stinner, A. Suárez, M. Winkler, D. Wrozek, A. Stevens, Othmer, D. Horstmann, A.Kubo, M. Peletier, I. Guerra, E. Espejo, C. Conca, Dolbeault, H. Zaag, Semba, Mizoguchi, T. Tsujikawa, Okuda, Calvez, Litcanu, Wolanski, J.I. Díaz Levine, Sleeman, M. Nilsen-Hamilton, T. Hillen, Gajewsky, Zacharias, B. Hu, M. Delgado, I. Gayte, Painter, Y. Tao, M. Negreanu, Wang, Cui, Izuhara. M. Fontelos, M. Vela etc

Dictyostelium discoideum (soil-living amoeba)

[Immune system](#) (white blood cells)

- Macrophages cells
- Neutrophil (granulocyte)

[Morphogenesis](#): the process of formation of the embryo.

See [Merkin-Needham-Sleeman](#) [2005], [Bollenbach et al](#) [2007],

[Angiogenesis](#)

- **Growth of tumours**. See [Anderson and Chaplain](#) [1998]
- Formation of embryos
- Healing of skin wound.

[Astrophysics](#) and gravitational interaction of particles.

See [Biler](#) [1995], [Biler-Hilhorst-Nadzieja](#) [1994]

2 Two species chemotaxis systems

We consider a two species system

$$u_t = \underbrace{d_1 \Delta u}_{\text{diffusion}} - \underbrace{\chi_1 \nabla \cdot (u \nabla w)}_{\text{chemotaxis}} + \underbrace{\mu_1 u (1 - u - a_1 v)}_{\text{proliferation and competition}}$$

$$v_t = \underbrace{d_2 \Delta v}_{\text{diffusion}} - \underbrace{\chi_2 \nabla \cdot (v \nabla w)}_{\text{chemotaxis}} + \underbrace{\mu_2 v (1 - a_2 u - v)}_{\text{proliferation and competition}}$$

$$\epsilon w_t = \underbrace{d_w \Delta w}_{\text{diffusion}} - \underbrace{\lambda w}_{\text{degradation}} + \underbrace{g(u, v)}_{\text{production}}$$

Neumann boundary conditions and appropriate initial data in Ω .

$$\epsilon = 0, \quad g(u, v) = u + v, \quad 0 \leq a_i < 1 \quad \text{for } i = 1, 2$$

We assume:

$$2(\chi_1 + \chi_2) + a_1 \mu_2 < \mu_1 \quad \text{and} \quad 2(\chi_1 + \chi_2) + a_2 \mu_1 < \mu_2.$$

The limit case

- $\chi_1 = \chi_2 = 0$ is a competitive system of two equations well studied in the literature. See [Pao \[1981\]](#), [Cosner and Lazer \[1984\]](#)

$$u_t = d_1 \Delta u + u(e_1 - b_1 u - c_1 v)$$

$$v_t = d_2 \Delta v + v(e_2 - b_2 u - c_2 v)$$

+ Neumann boundary conditions. Under assumption

$$\frac{e_1}{c_1} > \frac{e_2}{c_2}, \quad \frac{e_2}{b_2} > \frac{e_1}{b_1}$$

$$u \longrightarrow \frac{e_1 c_2 - c_1 e_2}{b_1 c_2 - c_1 b_2}, \quad v \longrightarrow \frac{b_1 e_2 - e_1 b_2}{b_1 c_2 - c_1 b_2}.$$

In our case, the assumption is equivalent to $0 < a_i < 1$ ($i=1,2$).

- $\mu_1 = \mu_2 = 0$. [Espejo-Arenas, Stevens and Velázquez \[2009\]](#)-[\[2010\]](#)
Simultaneous and non-simultaneous [blow up](#) of both species depending on the parameters and initial mass. See also [Espejo-Arenas and Conca \[2012\]](#), [Biler-Espejo-Guerra \[2013\]](#).

Steady states

$$\left\{ \begin{array}{l} 0 = \Delta u - \chi_1 \nabla \cdot (u \nabla w) + \mu_1 u (1 - u - a_1 v) \quad \text{in } \Omega, \\ 0 = \Delta v - \chi_2 \nabla \cdot (v \nabla w) + \mu_2 v (1 - a_2 u - v) \quad \text{in } \Omega, \\ 0 = \Delta w - \lambda w + u + v \quad \text{in } \Omega, \\ \frac{\partial u}{\partial n} = \frac{\partial v}{\partial n} = \frac{\partial w}{\partial n} = 0, \quad \text{in } \partial\Omega \end{array} \right.$$

The unique positive and bounded steady states is given by

$$u^* \equiv \frac{1 - a_1}{1 - a_1 a_2}, \quad v^* \equiv \frac{1 - a_2}{1 - a_1 a_2}.$$

An Auxiliary System of ODEs

$$\begin{cases} \bar{u}' = \bar{u} [\mu_1 - (\mu_1 - \chi_1)\bar{u} - \chi_1\underline{u} + \chi_1\bar{v} - (\chi_1 + \mu_1 a_1)\underline{v}], & t > 0, \\ \underline{u}' = \underline{u} [\mu_1 - \chi_1\bar{u} - (\mu_1 - \chi_1)\underline{u} - (\chi_1 + \mu_1 a_1)\bar{v} + \chi_1\underline{v}], & t > 0, \\ \bar{v}' = \bar{v} [\mu_2 + \chi_2\bar{u} - (\chi_2 + \mu_2 a_2)\underline{u} - (\mu_2 - \chi_2)\bar{v} - \chi_2\underline{v}], & t > 0, \\ \underline{v}' = \underline{v} [\mu_2 - (\chi_2 + \mu_2 a_2)\bar{u} + \chi_2\underline{u} - \chi_2\bar{v} - (\mu_2 - \chi_2)\underline{v}], & t > 0, \end{cases}$$

with initial conditions

$$\bar{u}(0) = \bar{u}_0, \quad \underline{u}(0) = \underline{u}_0, \quad \bar{v}(0) = \bar{v}_0, \quad \text{and} \quad \underline{v}(0) = \underline{v}_0.$$

The steady states of the ODEs system are given by

$$\bar{u}^* = \underline{u}^* = u^* \equiv \frac{1 - a_1}{1 - a_1 a_2}, \quad \bar{v}^* = \underline{v}^* = v^* \equiv \frac{1 - a_2}{1 - a_1 a_2}$$

We analyze the system of ODEs. under assumptions

$$0 < \underline{u}_0 < u^* < \bar{u}_0 < \infty, \quad 0 < \underline{v}_0 < v^* < \bar{v}_0 < \infty,$$

Step 1a.- $0 < \underline{u} < \bar{u} < \infty$ for $t \in (0, \infty)$;

Step 1b.- $0 < \underline{v} < \bar{v} < \infty$ for $t \in (0, \infty)$;

Step 2a.- $0 < \underline{u} < u^* < \bar{u} < \infty$ for $t \in (0, \infty)$;

Step 2b.- $0 < \underline{v} < v^* < \bar{v} < \infty$ for $t \in (0, \infty)$;

Step 3.- $\lim_{t \rightarrow \infty} |\underline{u} - \bar{u}| + |\underline{v} - \bar{v}| \rightarrow 0$.

Idea of Step 3.

$$\frac{d}{dt} \log \frac{\bar{u}}{\underline{u}} = \frac{\bar{u}_t}{\bar{u}} - \frac{\underline{u}_t}{\underline{u}} = -(\mu_1 - 2\chi_1)(\bar{u} - \underline{u}) + (2\chi_1 + \mu_1 a_1)(\bar{v} - \underline{v})$$

and

$$\frac{d}{dt} \log \frac{\bar{v}}{\underline{v}} = (2\chi_2 + \mu_2 a_2)(\bar{u} - \underline{u}) - (\mu_2 - 2\chi_2)(\bar{v} - \underline{v}) \quad \text{for all } t > 0.$$

We add both to obtain

$$\frac{d}{dt} \left(\log \frac{\bar{u}}{\underline{u}} + \log \frac{\bar{v}}{\underline{v}} \right) = (-\mu_1 + 2(\chi_1 + \chi_2) + \mu_2 a_2)(\bar{u} - \underline{u}) +$$

$$(-\mu_2 + 2(\chi_1 + \chi_2) + \mu_1 a_1)(\bar{v} - \underline{v})$$

for

$$\varepsilon := \min \{ \mu_1 - 2(\chi_1 + \chi_2) - \mu_2 a_2, \mu_2 - 2(\chi_1 + \chi_2) - \mu_1 a_1 \}$$

then

$$\frac{d}{dt} \left(\log \frac{\bar{u}}{\underline{u}} + \log \frac{\bar{v}}{\underline{v}} \right) \leq -\varepsilon (\bar{u} - \underline{u}) - \varepsilon (\bar{v} - \underline{v}) \quad \text{for all } t > 0,$$

This first entails that

$$\log \frac{\bar{u}}{\underline{u}} \leq \log \frac{\bar{u}_0}{\underline{u}_0} + \log \frac{\bar{v}_0}{\underline{v}_0} := c_0 \quad \text{for all } t > 0,$$

and

$$\log \frac{\bar{u}^*}{\underline{u}} \leq c_0 \implies \underline{u} \geq \bar{u}^* e^{-c_0} \geq \bar{u}^* \frac{\underline{u}_0}{\bar{u}_0} > 0 \quad \text{for all } t > 0.$$

In the same way

$$\underline{v} \geq \bar{v}^* \frac{u_0}{u_0} \frac{v_0}{v_0} > 0 \quad \text{for all } t > 0.$$

By the Mean Value Theorem

$$\bar{u}(t) - \underline{u}(t) = e^{\xi_1(t)} (\log \bar{u}(t) - \log \underline{u}(t)) \quad \bar{v}(t) - \underline{v}(t) = e^{\xi_2(t)} (\log \bar{v}(t) - \log \underline{v}(t)).$$

and

$$\frac{d}{dt} \left(\log \frac{\bar{u}}{\underline{u}} + \log \frac{\bar{v}}{\underline{v}} \right) \leq -\varepsilon_0 \left(\log \frac{\bar{u}}{\underline{u}} + \log \frac{\bar{v}}{\underline{v}} \right) \quad \text{for all } t > 0$$

is valid with

$$\varepsilon_0 = \varepsilon \frac{u_0}{u_0} \frac{v_0}{v_0} \min\{\bar{u}^*, \bar{v}^*\}.$$

After routine computations

$$0 < \log \frac{\bar{u}}{\underline{u}} \leq e^{-\varepsilon_0 t} \log \frac{\bar{u}_0 \bar{v}_0}{\underline{u}_0 \underline{v}_0}, \quad 0 < \log \frac{\bar{v}}{\underline{v}} \leq e^{-\varepsilon_0 t} \log \frac{\bar{u}_0 \bar{v}_0}{\underline{u}_0 \underline{v}_0} \quad \text{for all } t > 0$$

and thereby shows that

$$|\bar{u}(t) - \underline{u}(t)| + |\bar{v}(t) - \underline{v}(t)| \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

which ends the analysis of the ODE system.

Comparison: PDE system- ODE system
Reaction Diffusion Systems: [Rectangle Method](#) see [Pao \[1981\]](#).

We prove that, under assumption

$$0 < \underline{u}_0 \leq u_0(x) \leq \bar{u}_0 \quad \text{and} \quad 0 < \underline{v}_0 \leq v_0(x) \leq \bar{v}_0 \quad \text{for all } x \in \Omega.$$

we have

$$\underline{u} \leq u \leq \bar{u} \quad \text{and} \quad \underline{v} \leq v \leq \bar{v}.$$

We construct the functions:

$$\begin{aligned} \bar{U}(x, t) &:= u(x, t) - \bar{u}(t), & \underline{U}(x, t) &:= u(x, t) - \underline{u}(t), \\ \bar{V}(x, t) &:= v(x, t) - \bar{v}(t), & \underline{V}(x, t) &:= v(x, t) - \underline{v}(t) \end{aligned}$$

which satisfy

$$\begin{aligned} \bar{U}_t - \Delta \bar{U} + \chi_1 \nabla \bar{U} \nabla w &= \bar{U} [\mu_1 + (\chi_1 - \mu_1)(u + \bar{u}) + (\chi_1 - \mu_1 a_1)v - \chi_1 \lambda w] \\ &\quad + \chi_1 \bar{u} \bar{V} - \mu_1 a_1 \bar{u} \underline{V} + \chi_1(\underline{u} + \underline{v} - \lambda w). \end{aligned}$$

we multiply by \bar{U}_+ and after integration

$$\begin{aligned} \frac{d}{dt} \frac{1}{2} \int_{\Omega} \bar{U}_+^2 + \int_{\Omega} |\nabla \bar{U}_+|^2 &= -\frac{\chi_1}{2} \int_{\Omega} \nabla \bar{U}_+^2 \nabla w + \int_{\Omega} b(x, t) \bar{U}_+^2 \\ &\quad + \chi_1 \int_{\Omega} \bar{u} \bar{V} \bar{U}_+ - \mu_1 a_1 \int_{\Omega} \bar{u} \underline{V} \bar{U}_+ + \chi_1 \int_{\Omega} (\underline{u} + \underline{v} - \lambda w) \bar{U}_+ \end{aligned}$$

we obtain similar expressions for \underline{U} , \bar{V} , \underline{V} .

After many routine computations we obtain, for any $t < T$

$$\frac{d}{dt} \int_{\Omega} (\bar{U}_+^2 + \underline{U}_-^2 + \bar{V}_+^2 + \underline{V}_-^2) \leq k(T) \int_{\Omega} (\bar{U}_+^2 + \underline{U}_-^2 + \bar{V}_+^2 + \underline{V}_-^2)$$

Gronwall's lemma gives

$$\bar{U}_+ = \underline{U}_- = \bar{V}_+ = \underline{V}_- = 0.$$

3.- Competitive Exclusion (Stinner-T-winkler 2013)

We consider the system

$$u_t = \underbrace{d_1 \Delta u}_{\text{diffusion}} - \underbrace{\chi_1 \nabla \cdot (u \nabla w)}_{\text{chemotaxis}} + \underbrace{\mu_1 u (1 - u - a_1 v)}_{\text{proliferation and competition}}$$

$$v_t = \underbrace{d_2 \Delta v}_{\text{diffusion}} - \underbrace{\chi_2 \nabla \cdot (v \nabla w)}_{\text{chemotaxis}} + \underbrace{\mu_2 v (1 - a_2 u - v)}_{\text{proliferation and competition}}$$

$$0 = \underbrace{d_w \Delta w}_{\text{diffusion}} - \underbrace{\lambda w}_{\text{degradation}} + \underbrace{ku + v}_{\text{production}}$$

Neumann boundary conditions and appropriate initial data in Ω .

We define the new parameters

$$q_1 := \frac{\chi_1}{\mu_1} \quad \text{and} \quad q_2 := \frac{\chi_2}{\mu_2}.$$

We consider the assumptions

$$a_1 > 1 > a_2$$

k, q_1 and q_2 are nonnegative and $q_1 = \frac{\chi_1}{\mu_1} \leq a_1, q_2 = \frac{\chi_2}{\mu_2} < \frac{1}{2}$ and

$$kq_1 + \max \left\{ q_2, \frac{a_2 - a_2q_2}{1 - 2q_2}, \frac{kq_2 - a_2q_2}{1 - 2q_2} \right\} < 1.$$

The assumptions are equivalent to $kq_1 + q_2 < 1$ and

$$\begin{cases} kq_1 + (2 - a_2)q_2 + a_2 - 2kq_1q_2 < 1 & \text{if } kq_2 < a_2, \\ kq_1 + (2 - a_2 + k)q_2 - 2kq_1q_2 < 1 & \text{if } kq_2 \geq a_2. \end{cases}$$

Then, we have that

$$u \longrightarrow 0, \quad v \longrightarrow 1$$

Notice that the assumptions in the prototypical case $\chi_1 = \chi_2$, $\mu_1 = \mu_2$ are reduced to

$$\frac{\chi}{\mu} < \begin{cases} \frac{2+k-a_2-\sqrt{(k+2-a_2)^2-8k(1-a_2)}}{4k} & \text{if } a_2 > kq \\ \frac{2+2k-a_2-\sqrt{(2k+2-a_2)^2-8k}}{4k} & \text{if } a_2 \leq kq. \end{cases}$$

If we moreover have $k = 1$ then

$$\frac{\chi}{\mu} < \begin{cases} \frac{4-a_2-\sqrt{8-8a_2+a_2^2}}{4} & \text{if } a_2 \leq q \\ \frac{1-a_2}{2} & \text{if } a_2 > q. \end{cases}$$

In the limit case $k = 0$,

$$\frac{\chi}{\mu} < \frac{1-a_2}{2-a_2}$$

The borderline case $a_2 = 0$ reads

$$\frac{\chi}{\mu} < \frac{1}{2}$$

already found in [T-winkler \[2007\]](#)

Non-local terms (Negreanu-T [2013])

$$\begin{cases} u_t - \Delta u = -\chi_1 \nabla \cdot (u \nabla w) + u (a_0 - a_1 u - a_2 v - a_3 \int_{\Omega} u - a_4 \int_{\Omega} v), \\ v_t - \Delta v = -\chi_2 \nabla \cdot (v \nabla w) + v (b_0 - b_1 u - b_2 v - b_3 \int_{\Omega} u - b_4 \int_{\Omega} v), \\ -\Delta w + \lambda w = f + k_1 u + k_2 v, \end{cases}$$

with the homogeneous Neumann boundary conditions

$$\frac{\partial u}{\partial \nu} = \frac{\partial v}{\partial \nu} = \frac{\partial w}{\partial \nu} = 0, \quad x \in \partial\Omega, \quad t > 0,$$

and initial data

$$u(x, 0) = u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega.$$

“Global Competition”

Members of one species u compete for a limited resource z satisfying

$$u_t - \Delta u + \mu u = zu, \quad x \in \Omega, \quad t > 0$$

KPP-Fisher equation $z = (1 - u)$.

We assume that the resources diffuse and degrade with large diffusion coefficient ϵ^{-1}

$$-\frac{1}{\epsilon} \Delta z_\epsilon + \alpha_1 z_\epsilon = 1 - \alpha_2 u \quad x \in \Omega$$

$$\frac{\partial z_\epsilon}{\partial \vec{n}} = 0.$$

After Integration we have

$$\int_\Omega z_\epsilon = \frac{|\Omega|}{\alpha_1} - \frac{\alpha_2}{\alpha_1} \int_\Omega u.$$

Multiplying by u and after integration and thanks to Young inequality we have

$$\int_{\Omega} |\nabla z_{\epsilon}|^2 dx \leq \epsilon(1 + c(\alpha_1, \alpha_2)) \int_{\Omega} u^2 \rightarrow 0 \quad \text{as } \epsilon \rightarrow 0.$$

Then

$$z_{\epsilon} \rightarrow \text{constant} := \frac{|\Omega|}{\alpha_1} - \frac{\alpha_2}{\alpha_1} \int_{\Omega} u.$$

the equation

$$u_t - \Delta u + \mu u = zu, \quad x \in \Omega, \quad t > 0$$

is replaced by

$$u_t - \Delta u + \mu u = \mu u(1 - a_3 \int_{\Omega} u), \quad x \in \Omega, \quad t > 0.$$

If $\alpha_2 = \alpha_2(x)$ then, the nonlocal term $\int_{\Omega} \alpha_3(x)u$.

For the two species chemotaxis system we consider

$$-\frac{1}{\epsilon} \Delta z_{\epsilon} + \alpha_1 z_{\epsilon} = 1 - \alpha_2 u - \alpha_3 v \quad x \in \Omega \quad + \quad NBC$$

$$z_{\epsilon} \rightarrow \alpha - a_3 \int_{\Omega} u - a_4 \int_{\Omega} v.$$

Under assumptions

$$\int_0^\infty \left| \sup_{x \in \Omega} f - \inf_{x \in \Omega} f \right| \leq C_0 < \infty.$$

$$\chi_1, \chi_2, k_1, k_2, a_i, b_i > 0, \text{ for } i = 1, 2,$$

$$a_i \in \mathbb{R}, b_i \in \mathbb{R}, \text{ for } i = 3, 4$$

$$a_1 > 2k_1(\chi_1 + \chi_2) + b_1 + |b_3| + |a_3| \text{ and } b_2 > 2k_2(\chi_1 + \chi_2) + a_1 + |a_4| + |b_4|$$

we obtain the asymptotic behavior for positive initial data

$$u(\cdot, t) \longrightarrow u^* \equiv \frac{a_0(b_2 + b_4) - b_0(a_2 + a_4)}{(b_2 + b_4)(a_1 + a_3) - (b_1 + b_3)(a_2 + a_4)}$$

$$v(\cdot, t) \longrightarrow v^* \equiv \frac{a_0(b_1 + b_3) - b_0(a_1 + a_3)}{(b_1 + b_3)(a_2 + a_4) - (b_2 + b_4)(a_1 + a_3)}.$$

Two species Parabolic-ODE chemotactic system

$$u_t = \underbrace{d_1 \Delta u}_{\text{diffusion}} - \underbrace{\nabla(u\chi_1(w))}_{\text{chemotaxis}}$$

$$v_t = \underbrace{d_2 \Delta v}_{\text{diffusion}} - \underbrace{\nabla(v\chi_2(w))}_{\text{chemotaxis}}$$

$$w_t = h(u, v, w)$$

We assume

$$\chi_i, h \in W_{loc}^{1,\infty}(\mathbb{R}_+^2 \times \mathbb{R}), \quad \chi_i > 0.$$

$$\frac{\partial h}{\partial u} \geq \epsilon_u > 0 \quad \text{and} \quad \frac{\partial h}{\partial v} \geq \epsilon_v > 0$$

$$\frac{\partial h}{\partial w} < 0.$$

There exists w^* such that

$$h(u^*, v^*, w^*) = 0$$

where

$$u^* = \frac{1}{|\Omega|} \int_{\Omega} u_0 dx, \quad v^* = \frac{1}{|\Omega|} \int_{\Omega} v_0 dx.$$

Consequently (u^*, v^*, w^*) is a constant stationary solution of the system.

Global existence of solutions.

We assume

$$-h(0, 0, w) \leq \frac{k_i}{\chi_i(w)} \text{ for some } k_i > 0,$$

$$0 < k_{0i} \leq \chi_i(w) e^{\int_{\underline{w}}^w \chi_i(s) ds} \text{ for } w > \underline{w},$$

for $k_{0i} > 0$, with $i = 1, 2$. There exists \bar{u} and \bar{v} such that

$$h(u, v, \underline{w}) \geq 0, \quad h(u, v, \bar{w}) \leq 0, \quad \text{for } 0 \leq u \leq \bar{u}, \quad 0 \leq v \leq \bar{v},$$

where

$$\bar{u} := f_1(\bar{w}) \max \left\{ k_1 (\epsilon_u k_{01})^{-1}, \|u_0\|_{L^\infty(\Omega)} \right\},$$

$$\bar{v} := f_2(\bar{w}) \max \left\{ k_2 (\epsilon_v k_{02})^{-1}, \|v_0\|_{L^\infty(\Omega)} \right\},$$

for f_i defined by

$$f_i(w) = e^{\int_{\underline{w}}^w \chi_i(s) ds} \quad i = 1, 2.$$

Then, by using an [iterative method](#) ([Alikakos-Mosher iteration](#)) we have global boundedness.

Stability of the homogeneous steady states.

We consider

There exists $\alpha \in (0, 1)$ such that

$$\alpha h_w + u h_u \chi_1 + v h_v \chi_2 < 0 \quad \text{and} \quad 2\sqrt{1 - \alpha} h_w + u h_u \chi_1 + v h_v \chi_2 < 0.$$

Using a Lyapunov functional, we get that the steady state is globally asymptotically stable.

Notice that the previous assumptions are satisfied, for instance for

$$h(u, v, w) = u + v - w, \quad \chi_i(w) = \frac{\gamma_i}{1 + \gamma_i w} \quad \text{for } \gamma_i < 1/4$$

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Thank you very much for your attention!

